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**Performance of an Optically Read-Out Time Projection Chamber  
with ultra-relativistic electrons**

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# Performance of an Optically Read-Out Time Projection Chamber with ultra-relativistic electrons

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## Abstract

The Time Projection Chamber (TPC) is an ideal candidate to finely study the charged particle ionization in a gaseous medium. Large volumes TPCs can be readout with a suitable number of channels offering a complete 3D reconstruction of an ultra-relativistic charged particle track, that is the sequence of its energy releases in the TPC gas volume. Moreover, He-based TPCs are very promising to study keV energy particles as nuclear recoils, opening the possibility for directional searches of Dark Matter (DM) and the study of Solar Neutrinos (SN).

In this paper we report the analysis of the data acquired with a small TPC prototype (named LEMOn) built by the CYGNO collaboration that was ex-

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posed to a beam of 450 MeV electrons at the Beam Test Facility of National Laboratories of Frascati. LEMOn is operated with a He-CF<sub>4</sub> mixture at atmospheric pressure and is based on a Gas Electron Multipliers amplification stage that produces visible light collected by a sub-millimeter position resolution scientific CMOS camera. This type of readout - in conjunction with a fast light detection - allows a 3D reconstruction of the electrons tracks. The electrons are leaving a trail of segments of ionizations corresponding to a few keV energy release each. Their study leads to predict a keV energy threshold and 1-10 mm longitudinal and 0.1-0.3 mm transverse position resolution for nuclear recoils, very promising for the application of optically readout TPC to DM searches and SN measurements.

*Keywords:* TPC, MPGD, GEM, Optical Read Out, Tracking detector, Dark Matter

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## 1. Introduction

Large Time Projection Chambers (TPC) have been employed in various high energy physics and nuclear physics experiments. These detectors consist of a relatively large gas volume to track a charged particle's trajectory, that is measuring the position of the ionization trail of the charged particle in the gas with sub-mm spatial resolution. TPC are among the best devices when used in a magnetic field to measure the momentum of a relativistic charged particle and to allow its mass identification by detecting its specific ionization. The absolute distance from the particle's trajectory to a readout plane (the *longitudinal* coordinate) is usually obtained by measuring the drift time of the ionization electrons in a motion along a uniform electric field in the gas. The *transverse* coordinates are usually obtained with a read-out plane finely segmented with a wire, strip or pixel geometry. Finally, an electronic signal is formed in an amplification stage that is producing an avalanche out of the initial ionization electrons.

TPC have been also proposed for ultra-rare events searches as nuclear and

electron recoils in the directional search for Dark Matter (DM) and results has been recently reported on these searches [1, 2, 3]. Different types of electron amplification and readout have been studied[4] based on a multi-wire proportional chamber scheme (DRIFT[2, 5]), on micro-pattern gaseous detectors (MPGD) using pixels or strips and a custom electronics for a complete 3D readout (MIMAC[6, 7], NEWAGE [8], D3 [9]). Also the detection of neutrinos coming from the Sun (SN) with large TPC was proposed [10, 11].

An alternative read-out technique (referred as *optical read-out*) is based on the detection of the light emitted by the gas molecules during the avalanche processes triggered by the primary ionization electrons[12, 13, 14]. In this approach, the use of suitable lenses allows the reconstruction of the projection of the particle trajectory on the readout plane.

A first directional DM detector prototype was using an optical read-out with an image intensified CCD [15]. The DMTPC [16] collaboration was then proposing to use a CCD with a wire mesh amplification stage in a low pressure  $\text{CF}_4$  gas. A combined readout of a PMT and CCD was also attempted for a 3D reconstruction [17]. The use of Gas Electron Multipliers (GEM) [18] as MPGD amplifiers was then introduced to study low energy electron and nuclear recoils in low pressure  $\text{CF}_4$  with a CCD readout [19]. Moreover, this optical approach was tested and developed in past years for other applications [20, 21, 22].

The recent introduction of low noise and high granularity scientific CMOS (sCMOS) sensors represented an important step forward in terms of granularity and readout noise with respect to the use of CCD. The better performance reflects on a better space resolution and a higher sensitivity that allows to clearly detect and reconstruct very low energy releases while working at atmospheric pressure [23, 24].

The MPGD optical readout can in fact represent an opportunity for large TPCs for rare events searches. Optical sensors offer higher granularity with respect to electron sensitive devices and the optical coupling locates the sensors out of the sensitive volume, reducing the interference with the MPGD high voltage operation and the gas contamination. Moreover, the use of suitable

lenses enables to image large surfaces to small sensors.

In last years, optical sensors market had an impressive development: sensors able to provide the required large granularity along with a very low noise level and high sensitivity to single photon counting are now available. On the other hand, the readout frame rate well below the kHz value still represents a limiting factor for these devices. The radiopurity of the optical sensors could also be an element to be improved for a future DM detector based on this technology.

Many smaller scale prototypes have been built and tested so far showing promising results (NITEC [25], ORANGE [26, 27]).

In this paper we describe the data acquired with a 7 litres prototype (named LEMOn - Long Elliptical MOdule) exposed to 450 MeV electrons at the Beam Test Facility (BTF) [28] of INFN National Laboratory of Frascati (LNF). LEMOn has an optical readout of a triple-GEM structure based on one sCMOS camera to image energy releases in He-based gas mixtures kept at atmospheric pressure. This would therefore allow to host larger target mass for DM searches than previous low pressure TPC prototypes and experiments.

BTF ultra-relativistic electrons are leaving along their path a straight trail of clusters of few ionization electrons each. We evaluate the capability to detect within the LEMOn drift region short sequences of ionization clusters (*segments* of the track) and to reconstruct their positions. Each of these segments corresponds to small energy deposits of the order of one keV. Even if the topology of the track segments is different from that typical of the nuclear or electron recoils of the same energy, with this method keV energy releases can be easily characterized. This would open the possibility to operate He-based gas mixture TPC with a keV energy threshold, to be sensitive to a DM particle with a mass in the range of GeV.

The use of a beam with a precise time reference allows to operate LEMOn as a standard TPC to evaluate the absolute longitudinal distance. However, in absence of a reference time for the events as in a DM search, the ionization electron drift time can not be exploited to determine their initial longitudinal position inside the detector. In this paper, we study also a method to measure

the longitudinal position in the drift volume based on the electron diffusion, as proposed for an electronic readout[29, 30] too. This is in fact very important to fully define a fiducial gas volume in a DM search, in a way to discard the background events generated by the detector components, in particular the cathode and the amplification system (as the GEMs)[31, 32].

LEMO<sub>n</sub> was also tested with other particles sources [33, 34, 35, 36] as neutron beams and various radioactive sources. A design of a 1 m<sup>3</sup> demonstrator (named CYGNO) to be constructed in 2021-2022 and to be hosted at the INFN National Laboratory of Gran Sasso (LNGS) will rely on the LEMO<sub>n</sub> performance. The main sizes of LEMO<sub>n</sub> (20 cm maximum drift path and almost 500 cm<sup>2</sup> GEMs) are in fact of the same order of magnitude of the typical sizes of readout modules foreseen for the 1 m<sup>3</sup> scale (50 cm maximum drift path and around 1000 cm<sup>2</sup> GEMs).

In a later phase, a 30-100 m<sup>3</sup> detector is foreseen, as an element of a world distributed observatory for DM and SN within the CYGNUS international network[37, 38, 39].

## 2. LEMO<sub>n</sub> Prototype Design

The LEMO<sub>n</sub> prototype structure (Fig. 1) was made of Acrylic Styrene Acrylonitrile (ASA) at the 3D printing Facility of LNF[40]. This has offered the opportunity to easily design and to quickly develop detectors and also to test the 3D printing system for gas detector applications.

The LEMO<sub>n</sub>'s heart consists of a 7 liter active drift volume surrounded by a 200×240 mm<sup>2</sup> elliptical field cage with a 200 mm distance between the anode and the cathode. The electric field is shaped by 1 mm Cu+Ag wires held at their positions by nineteen 3D-printed rings with 1 cm pitch. The anode side is instrumented with a 200×240 mm<sup>2</sup> rectangular triple GEM structure located at a position 10 mm apart from the last field cage ring. They are LHCb-like [41] 50 μm thick GEMs, with 70 μm diameter holes and 140 μm pitch and with two 2 mm wide transfer field gaps between them (Fig. 2).

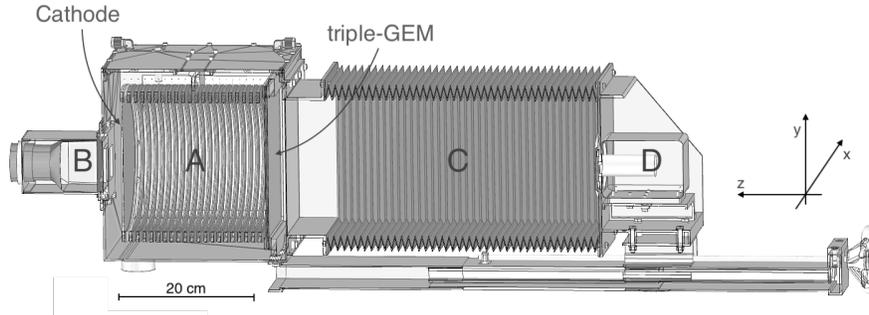


Figure 1: LEMOn design. PMT holder (B), semi-transparent cathode, field cage rings (A), the triple GEM stack, the optical bellow (C) and the ORCA Flash camera holder (D). Reference frame is also indicated: the origin of the axis is located at the middle of the GEM plane in vertical Y coordinate, while the X origin is located at the beginning of the GEM plane to coincide with the particle track length; moreover, the Z coordinate represents the distance from the GEM plane starting from it.

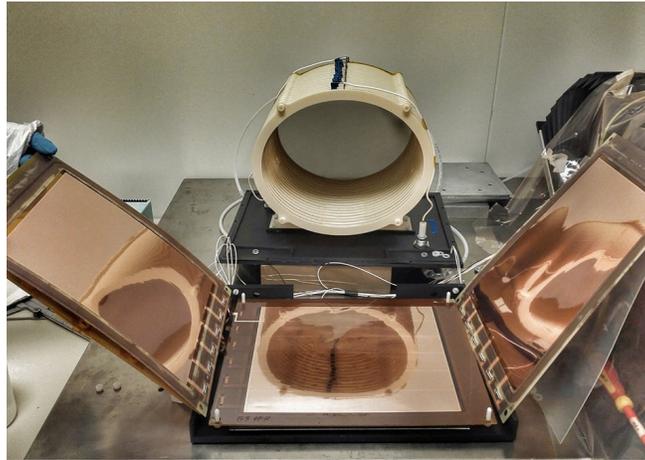


Figure 2: Exploded triple GEM structure (on the front) with the 3D-printed field cage rings (white) equipped with a semi-transparency cathode (on the back of the rings).

A  $203 \times 254 \times 1 \text{ mm}^3$  transparent window and a bellow with a tunable length allows to collect the light emitted within the GEMs by using an ORCA-Flash 4.0 camera [42] positioned at a distance of about 50 cm from the outermost GEM layer and based on a sCMOS sensor. The main feature of CMOS technology

is that each individual pixel has its own amplifier. Therefore, pixel responses can be quickly processed, suppressing possible noise pick-up [43]. Moreover, individual amplifiers allow to operate CMOS based sensors at a higher frame rate with respect to other technologies (e.g. CCD devices [44]). ORCA-Flash 4.0 camera is provided with a high granularity ( $2048 \times 2048$  pixels), very low noise (standard deviation of readout noise around two photons per pixel and negligible effect of dark current for exposure times lower than 1 second), good sensitivity (70% quantum efficiency at 600 nm, comparable with other front-illuminated devices) and good linearity.

This camera is instrumented with a Schneider lens (25 mm focal length and aperture of  $f/0.95$ ). The lens is placed at a distance  $d=50.6$  cm from the last GEM in order to obtain a de-magnification  $\delta = (d/f) - 1 = 19.25$  to image a surface  $25.6 \times 25.6$  cm<sup>2</sup> onto the  $1.33 \times 1.33$  cm<sup>2</sup> sensor. In this configuration, each pixel is therefore imaging an effective area of  $125 \times 125$   $\mu\text{m}^2$  of the GEM layer. Because of the geometrical acceptance [23], the fraction of the light collected by the lens can be evaluated to be  $1.7 \times 10^{-4}$ . Except from very few pixel rows and columns at the sensor sides, that look outside from the TPC sensitive volume and are removed from the analysis, sensor response and noise level do not show any appreciable dependence on  $X$  and  $Y$ .

A cathode has been realized using an ATLAS MicroMegas mesh [45], produced by Swiss BOOP company with  $30$   $\mu\text{m}$  diameter wires with a pitch of  $70$   $\mu\text{m}$ , stretched and glued on a ring,  $1$  cm apart from the first ring of the field cage and ensuring an adequate light transmission. Its transparency to light has been estimated to be about 70%.

On this side, LEMOn has been equipped with a  $50 \times 50$  mm<sup>2</sup> HZC Photonics XP3392 photomultiplier [46] (PMT) detecting light through a transparent  $50 \times 50 \times 4$  mm<sup>3</sup> fused silica window.

This PMT also allowed to study the timing properties of light produced in the innermost GEM layer  $250$  mm far away during the development of the electron avalanche. The field cage is contained in a  $370 \times 270 \times 280$  mm<sup>3</sup> and  $2.5$  mm thick box made of ASA. This box, realised to guarantee the gas

containment, has been equipped with two  $180\ \mu\text{m}$  thin  $200 \times 200\ \text{mm}^2$  windows made of TEDLAR to reduce as much as possible the multiple scattering of the impinging ultra-relativistic electrons (or the absorption of particles in other tests).

### 3. Detector Operations

The data reported in this paper have been collected at the BTF which can deliver bunched electrons or positrons from few tens of MeV to several hundred MeV energy. The BTF is optimized to deliver electrons with an energy of 450 MeV with a typical beam spot size of  $\sigma_{x,y} \simeq 2\ \text{mm}$ , a divergence  $\sigma'_{x,y} \simeq 2\ \text{mrad}$ , and an energy spread of  $\simeq 1\%$ . Moreover, for our measurements the bunch multiplicity has been tuned to be single particle, with the possibility to change it up to several hundred electrons (with a 10 ns bunch length). The bunch multiplicity and the beam spot-size were monitored by means of a silicon pixel detector (Fast interface for Timepix pixel detectors, FitPix [47]) located upstream of LEMOn and a Pb-glass calorimeter [48, 49] accommodated at the back of LEMOn and acquired together with the LEMOn camera image and the PMT signal. All the devices in the acquisition chain were triggered by the DAPHNE accelerator timing system of which the BTF is part [50], that allows to distribute the trigger to the camera, to the FitPix, and to a LeCroy 610Zi used to digitise at a sampling frequency of 10 GS/s and acquire waveforms from the LEMON-PMT and from the calorimeter.

LEMON was operated with a He-CF<sub>4</sub> (60/40) gas mixture, with a voltage of 455 V across the sides of each of the three GEMs and an electric field between them of 2.0 kV/cm. In this high voltage configuration the GEM stack provides a gain of about  $5 \times 10^5$ , considered large enough to reach a good light production in the avalanche generated in the GEMs structure. The typical photon yield for this type of gas mixtures has been measured to be around 0.07 photons per avalanche electron [23, 51, 52]. The gas mixture was kept at atmospheric pressure under continuous flow of about 200 cc/min.

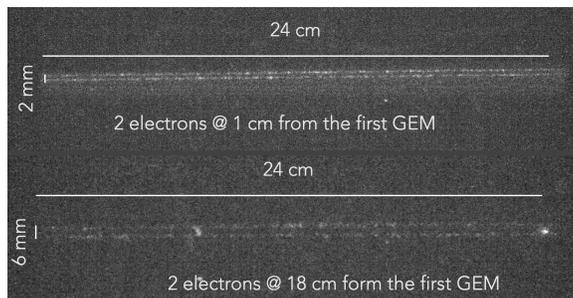


Figure 3: Two examples of two-tracks events collected at the Frascati BTF. The two BTF electrons are crossing the field cage parallel to the 240 mm major axis. One event is acquired with the beam close to the GEM (up), and another with the beam 180 mm far away from the GEM (down). The separation among the two track can be easily measured.

The field cage was powered by a High Voltage Power Supply CAEN N1570[53] able to deliver up to 15 kV. This posed a limitation to the maximum electric field at 0.6 kV/cm.

The triple GEM system was powered with a HV GEM power supply [54] ensuring stability and accurate monitoring of the bias currents.

The ORCA Camera I/O was configured to work in *Global Reset Mode*. The pixel lines are initialized one at a time. When the sensor is fully initialised, all previously collected charges are reset and all pixels are simultaneously exposed for the Global Exposure Time, configured by software.

The initialisation process lasts about 80  $\mu$ s and was triggered by a signal properly generated by the accelerator timing system previously described.

Optics and exposure time (30 ms) were optimized to ensure the largest light collection and to reduce as much as possible events due to the natural radioactivity. One hundred images were typically acquired in each detector configuration (*run* in the following). Fig. 3 shows two examples of BTF electron tracks images acquired with LEMOn.

The 450 MeV electrons were delivered by the LNF accelerator complex along the X axis direction (see Fig.1) with a repetition rate of one bunch every second. LEMOn was accommodated over a remotely controlled table in order to scan

the  $Z$  coordinate (with a 0.2 mm precision).

#### 4. Test Beam Results

Data were collected during a week long campaign, when LEMOn was continuously operated.

The fast PMT signal waveform was acquired using as external trigger from the timing signal of the BTF line synchronized with the electrons arrivals. The time  $t_s$  corresponding to a fixed voltage of the PMT waveform was associated to each PMT signal and data at different longitudinal  $Z$  position were collected. The standard deviation ( $\sigma_t$ ) of the distribution of the residuals of  $t_s$  at each  $Z$  can be converted into the standard deviation ( $\sigma_Z$ ) of the BTF electrons  $Z$  position using the gas mixture drift velocity. A value of  $\sigma_Z$  around 1 mm was found, well compatible with the beam spot transverse size.

Each image acquired with the sCMOS camera was saved as a 2048 x 2048 matrix of photon counts. Because of the very low occupancy of sensor, the baseline noise of the sensor is estimated pixel by pixel by obtaining the distribution of counts for each pixel in all the images of a data-taking run. The average count and the standard deviation ( $\sigma_n$ ) for each pixel is then evaluated. This average photon count is subtracted to the count of each image before the image is further processed[55].

The reconstruction of tracks in each image is then made by using a Hough transform pattern recognition algorithm (HT) [56]. Since LEMOn is positioned to let the BTF electrons cross the drift volume at  $Y \sim 0$ , only pixels with a photon count exceeding  $1.5\sigma_n$  are used in the HT. The HT could finds several lines connecting all the pixels above this threshold: the most ranked line (i.e. the one that minimise the sum of residuals to points) within an angle of about  $\pm 5$  mrad respect to the  $X$  axis is then selected in each image and it represents the candidate reconstructed BTF electron (a *track*). Multiple track images are also analyzed and the HT is in fact able to distinguish tracks in events with multiplicity larger than one.

The efficiency to reconstruct an electron crossing the 24 cm wide drift region was estimated by selecting a sample of events with the calorimeter signal compatible with a single BTF electron deposit. In all the events considered the HT was always able to find a track in the sCMOS image. This is well compatible with a Garfield [57, 58] simulation of the gas mixture that yields an estimate for the average energy loss per single ionization of 42 eV and about 5 primary electrons produced per millimeter, resulting in a  $dE/dx$  of about 0.21 keV/mm for a  $\simeq 450$  MeV electron. This energy loss translates into three primary ionization  $e^-$  cluster per track mm. The sequence of several ionization clusters along the electron trajectory is a signature clear enough to detect each electron with the HT.

#### *4.1. Position and energy resolutions*

The total light yield distribution transverse to the track direction is obtained for each track. This distribution is fitted to a Gaussian function and its integral is proportional to the total energy loss of the track in the gas volume of the field cage (right top insert of Fig. 4).

In all analyses presented in this paper, tracks have been divided in 36 portions about 7 mm long.

We evaluate the performance of LEMOn in measuring their characteristics: position in space, light content, longitudinal and transverse light profile. An average energy loss of about 1.5 keV corresponds to each segment (with a production of 21 clusters for a total of almost 40 primary electrons per segment) allowing to study the performance of LEMOn to in the keV energy release range.

Among these 36 segments only 18 in the central region of the field cage ellipse are retained for subsequent analysis. This is due to the circumstance that instrumental effects such as electric field distortion and optical vignetting, have been observed in the external regions of sensitive volume of LEMOn.

In different runs the LEMOn position along  $Z$  is changed resulting in a different  $Z$  coordinate for the tracks. These runs are used to evaluate the uniformity versus the drift axis ( $Z$ ) of the TPC. The total light yield per track shows a

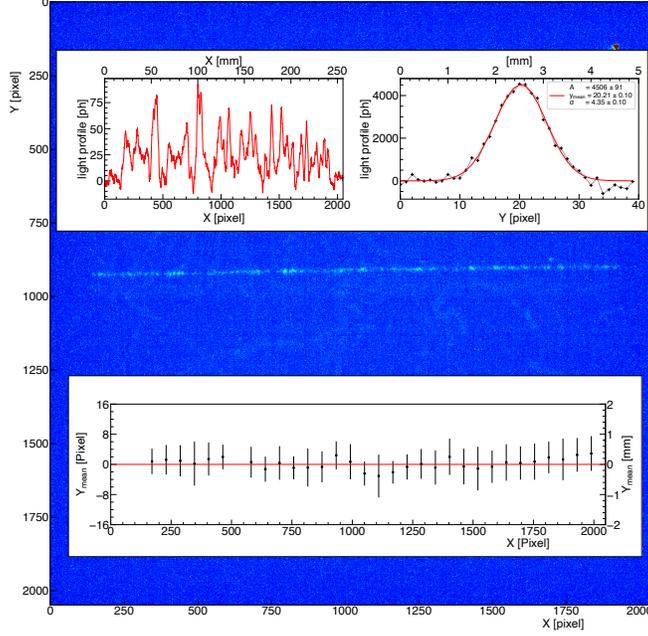


Figure 4: **Background:** sCMOS camera image collected at the BTF with beam passing 6.5 cm apart from GEM plane. Pixels collecting light due to a BTF electron are visible at  $Y \sim 1000$ . **Left top insert:** light distribution (after pixel-by-pixel noise subtraction) along the track  $X$ -direction summing all the photons for 40 pixel in the  $Y$  direction around the track direction. **Right top insert:** light distribution transverse to the track ( $Y$  direction) summing all the photons along the  $X$  direction with a superimposed Gaussian fit. **Bottom insert:** Gaussian  $Y_{mean}$  of the transverse light distribution with superimposed the line found by the HT (subtracted of the  $Y$  of the line). The error bar of the points is the sigma of the normal fit in each slice and it is taken as a indication of the width of the light deposit. Some segments with a too low signal-to-noise ratio of the detected light are not displayed.

decrease at larger  $Z$  positions (see Fig.5) resulting in a total light yield reduction of almost 40% for tracks 20 cm away from the GEMs. This is likely due to the electron attachment to gas impurities during the drift of the ionization electrons from their production points along the track to the GEMs. A very similar decrease of the light yield was also measured by the PMT, excluding this to be due to some instrumental or analysis effect. Moreover, the relative fluctuation of the total light yield per track represents an estimate of the energy release

resolution (Fig.6). This turns out to have a slight dependence on the  $Z$  position of the track increasing from about 20% at  $Z=0$  cm to about 30% at  $Z=20$  cm, corresponding to an absolute resolution between 300 eV and 450 eV. Since an average number of about 21 electron clusters per segment are produced, this result can almost completely be explained with the statistical fluctuations of gas ionization.

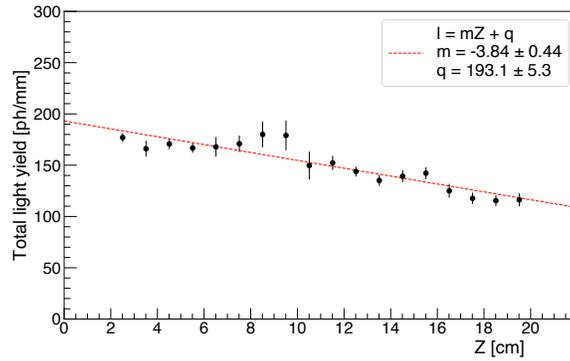


Figure 5: Average total light yield  $I$  per track as a function of the track  $Z$  position across the field cage.

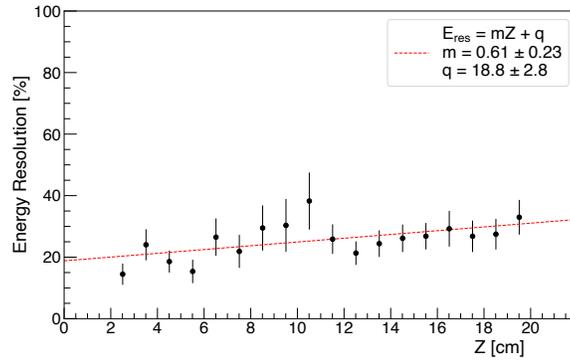


Figure 6: Energy Resolution: standard deviation of total light yield per track divided by its average as a function of the track  $Z$  position.

Also the transverse light profile of each segment can be described by a Gaus-

sian distribution. The fitted mean value  $Y_{mean}$  of this distribution is used to identify the segment's  $Y$  position. In Fig. 4 examples of longitudinal and transverse light profiles are shown. The distribution of the residuals of  $Y_{mean}$  to the line obtained with HT is then considered for the 18 segments. The standard deviation of this distribution evaluated in a sample of several tracks represents an estimate of the resolution on the segment's  $Y$  position. This procedure is repeated for several images acquired with the BTF electrons crossing the LEMOn field cage at different  $Z$  positions (Fig.7). The dependence of the segment's  $Y$  position resolution on the  $Z$  track coordinate is then interpolated with a linear function (Fig.7). A resolution of  $83 \pm 12 \mu\text{m}$  is extrapolated close to the GEMs ( $Z = 0$ ), in agreement with an almost three times better result obtained with an almost three times smaller effective pixel in a previous work [59]. We observe its worsening for larger  $Z$ : this is mainly due to the reduction of electrons, probably due to re-absorption, for larger  $Z$  that, besides reducing the produced light, also amplifies the effect of diffusion.

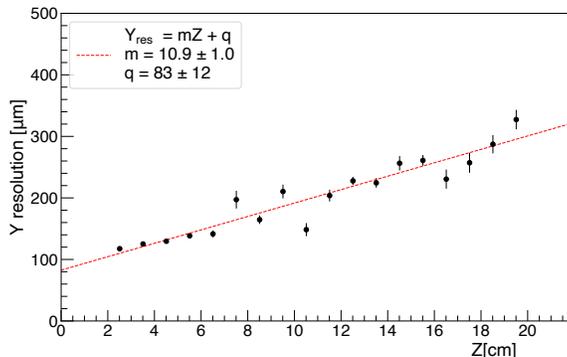


Figure 7: Segment  $Y$  position resolution as a function of the track's  $Z$  coordinate.

#### 4.2. $Z$ coordinate measurement

During the drift in the gas, electrons are subject to longitudinal and transverse diffusion effects, that influence their arrival  $X$  and  $Y$  coordinates and their arrival time at the GEMs. For each single electron, the position at the anode

has a probability distribution function that can be described by a 2D Gaussian with  $\sigma_X$  and  $\sigma_Y$  equal to  $B = \sqrt{\frac{2DZ}{\mu E}}$  where  $D$  is the diffusion coefficient,  $\mu$  the electron mobility and  $E$  the drift electric field [60]. Moreover, during the avalanche formation within the GEMs a further diffusion of the avalanche electrons is taking place. Eventually, the light recorded by the sCMOS camera and by the PMT is related to the original point with an uncertainty that is larger for production points farther in  $Z$  from the GEMs. This is reflected in the transverse ( $Y$ ) light distribution of each segment of the track: the  $\sigma_Y$  obtained from the Gaussian fit is in fact increasing with  $\sqrt{Z}$ . Therefore, the segment's original  $Z$  can be deduced by measuring  $\sigma$ (Fig.8).

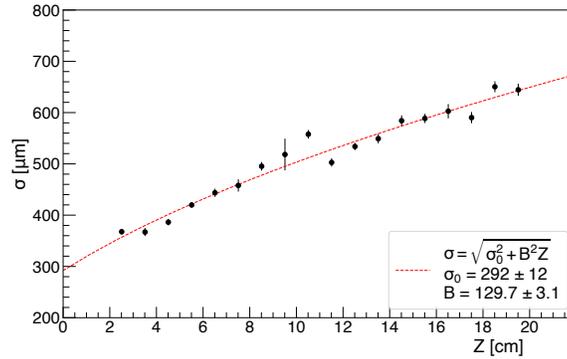


Figure 8: Average  $\sigma$  of transverse light distribution for track segments as a function of the track  $Z$  coordinate.

The observed values of  $\sigma$  are related to the track's  $Z$  by  $\sigma = \sqrt{\sigma_0^2 + B^2 Z}$

The transverse diffusion coefficient  $B$  in the gas is measured to be  $129.7 \pm 3.1 \frac{\mu m}{\sqrt{cm}}$ , well in agreement with the expected value obtained with Garfield simulation ( $130 \frac{\mu m}{\sqrt{cm}}$ ) [61].

The intercept at zero  $\sigma_0 = 292 \pm 12 \mu m$  is due to the contribution of the electron avalanche propagation in the GEM stack and confirms results obtained with radioactive source [61].

From the same Gaussian fit to the light  $Y$  distribution of each segment, the amplitude  $A$  can be obtained. Since  $\sigma A$  is proportional to the total light  $I$  of

the segment we can define  $\eta = \frac{\sigma}{A}$  that is therefore proportional to  $\frac{\sigma_0^2 + B^2 Z}{I}$ .

In Fig. 9 we show the dependence of  $\eta$  on the track's  $Z$  coordinate. Since in our data the light  $I$  shows a linear decrease with  $Z$ ,  $I(Z)$  can be described as  $I_0 - \epsilon Z$  (see Fig.5). For small values of  $\epsilon$ ,  $\eta \simeq (\sigma_0^2 + B^2 Z) \cdot (I_0 + \epsilon Z)$ . Therefore, a quadratic fit is expected to give a better representation of the data.

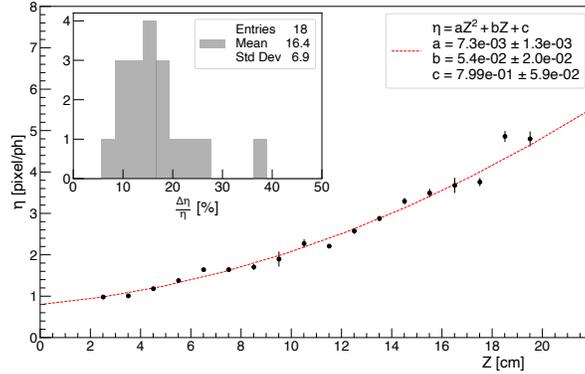


Figure 9: Average  $\eta$  for track segments as a function of the track's  $Z$  coordinate with a quadratic fit superimposed. Relative uncertainty on  $\eta$  (inset) for the various  $Z$  positions.

A similar parameter  $\eta_{PMT}$  can be defined from the analysis of the PMT waveform. In this case the total light of the track is recorded and we use the Pb-glass calorimeter signal of the BTF line to reject events with more than one track. The amplitude of the PMT waveform and its width can be similarly used to calculate  $\eta_{PMT}$ . In this case the width of the waveform is larger for more distant tracks (larger  $Z$ ) due to the longitudinal diffusion of the drifting ionization electrons. By using the drift velocity,  $\eta_{PMT}$  can be related to  $Z$  similarly to  $\eta$  (Fig.10).

Transverse and longitudinal diffusion can therefore be exploited to measure the longitudinal  $Z$  coordinates. We can estimate  $\frac{\Delta\eta}{\eta}$  from the standard deviation of the distributions of the  $\eta$  values. Since  $\frac{\Delta\eta}{\eta} = \frac{\Delta Z}{Z}$ , the latter can be used to estimate  $\frac{\Delta Z}{Z}$ . For each single segment, this resolution turns out to be in the range 10% - 20% for both methods, even if measurements with sCMOS lead to a slightly better evaluation of  $Z$ .

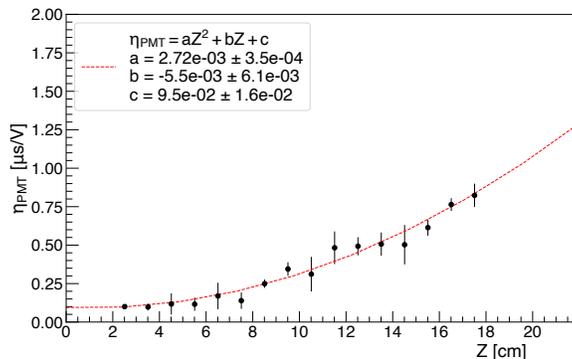


Figure 10: Average  $\eta_{PMT}$  as a function of the track  $Z$  coordinate. Data were not recorded for the largest  $Z$  point since the beam was adding extra noise in the PMT.

Moreover, a combination of the  $\eta$  parameters (from sCMOS and PMT) would probably give an even better evaluation  $Z$  that would result to be very useful to define a fiducial region of a DM detector. The anode where the GEMs are located (at  $Z \sim 0$  in LEMOn) and the cathode (at  $Z \sim 20$  cm in LEMOn) are usually the most radioactive elements in a DM detector. They are in fact sources of spurious nuclear recoils that would be easily removed if their  $Z$  coordinate is measured.

## 5. Conclusion

The data collected at the Frascati BTF with the LEMOn prototype, part of the CYGNO project, are confirming the potentiality of large optically readout TPC as detector for low energy events.

In this paper we have reported how ultra-relativistic electron tracks can be very efficiently reconstructed by collecting the light emitted from GEMs with a high resolution and high sensitivity sCMOS camera and a PMT. The analysis of the 7 mm long slices of an electron track shows a good spatial and energy resolution, making very promising the use of this gas TPC down to few keV energy releases. Also a method based on the ionization electron diffusion confirmed to be very effective in determining keV segment's longitudinal position

within the field cage.

The very good sensitivities measured in the low energy region make this technology extremely interesting in the development of a larger scale detector operated at atmospheric pressure aiming to observe very rare processes as Dark Matter or Solar Neutrinos interactions with ordinary matter.

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### **Data Availability Statement**

Data available on request from the authors.

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